Static and cyclic undrained response of square embedded foundations

N. NTRITSOS*, I. ANASTASOPOULOS† and G. GAZETAS‡

Results of a three-dimensional finite-element study for the effect of embedment on the undrained bearing capacity, the elastic stiffness, and the cyclic behaviour of square-in-plan foundations are presented. Uniaxial horizontal (Q) and pure-moment (M) limit loads, as well as the respective elastic stiffnesses (KHH and KMM) are obtained, and simplified models are developed to interpret the observed trends. The substantially different role of embedment in increasing elastic stiffness and in increasing ultimate loads is interpreted in simple soil mechanics terms. Extensive comparisons are made with the two-dimensional results for a strip foundation. Combined (QM) loading capacities are obtained and presented as ‘interaction’ diagrams; the significance of the vertical load (N) is also addressed. The importance of the type of contact between the foundation interfaces (vertical or horizontal) with the surrounding or underlying soil is explored. A substantial reduction in all capacities is shown when the interfaces are tensionless and of limited shear (sliding) resistance (TSI), compared with the capacities for the ideal case of fully bonded contact (FBC). The cyclic moment–rotation (M–θ) response of embedded foundations carrying a simple slender structure is investigated parametrically. It is found that the monotonic loading curves provide approximately the envelope for the cyclic M–θ loops. But the shape of these loops and the ensuing settlement of the foundation are both functions of the factor of safety (FSV) against vertical bearing capacity. In case of surface and shallowly embedded foundations with high values of FSV (i.e. with light loading or on very stiff soil) the loops pass nearly through the centre (M = θ = 0) of the coordinate axes and the residual settlement is negligible; both are a consequence of the predominantly geometric non-linearity in the form of separation of the walls and uplifting of the base from the soil. On the contrary, with low FSV values (heavy loading or soft soil) and deeper relative embedment, broad hysteresis loops and accumulating settlement are the rule – a product of material inelasticity of the soil, under limited soil–foundation detachment.

KEYWORDS: bearing capacity; failure; finite-element modelling; footings; foundations; soil/structure interaction; stiffness

INTRODUCTION

Static foundation design calls for avoiding failure of the soil–foundation system, and ensuring that the system displacements will be compatible with the desired structural performance. Prediction of the ultimate ‘failure’ loads of a foundation has traditionally been addressed in practice with the conventional bearing capacity theory (Terzaghi, 1943; Meyerhof, 1953; Brinch Hansen, 1970). The basis of this theory is the algebraic expression of the ultimate vertical load \( N_u \) as a function of the soil strength parameters, the soil mass density and the foundation geometry. When vertical, horizontal and moment loading act simultaneously on the foundation, to obtain the ultimate condition a series of ‘reduction’ factors for inclination and eccentricity are applied to \( N_u \). The effect of embedment is incorporated in the bearing capacity expression with the use of semi-empirical depth factors increasing the capacity over that of the surface footing with the same base geometry. Although this traditional approach has gained widespread acceptance in foundation engineering practice, its quasi-empirical nature renders it just an approximation. Moreover, it cannot be directly used in numerical one-step soil–structure interaction analyses, as explained by Houlby (2003). To overcome drawbacks, many researchers (Butterfield & Ticoff, 1979; Georgiadis & Butterfield, 1988; Nova & Montrassio, 1991; Butterfield & Gottardi, 1994; Martin, 1994; Salencon & Pecker, 1995; Bransby & Randolph, 1999; Gottardi et al., 1999; Houlby & Puzrin, 1999; Taiebat & Carter, 2000; Houlby, 2003; Randolph & Puzrin, 2003; Gourvenec, 2007, 2008; Yun & Bransby, 2007; Chatzigogos et al., 2009; Randolph & Gourvenec, 2011) have developed the concept of a ‘failure envelope’ which in a Cartesian coordinate system with axes the vertical force, \( N \), the horizontal force, \( Q \), and the overturning moment, \( M \), is the surface that indicates bearing capacity failure for all possible \( NQM \) load combinations. When the point representing a particular load combination plots inside this ‘failure envelope’, the foundation does not fail; when just touching the envelope, it mobilises a ‘failure’ mechanism and deforms inelastically.

On the other hand, seismic design of soil–foundation–structure systems requires knowledge of the elastic or quasi-elastic stiffnesses of the foundation, since these more-or-less control the natural frequencies of the system, and hence the loads that are transmitted from the structure to the foundation. A significant effort has been spent on obtaining static and dynamic stiffnesses of foundations; a wealth of results has been published in the last 40 or more years, for both uniaxial and combined loading on various idealised soil profiles: compilations of such solutions have appeared in papers by Roesset (1980), Gazetas (1983), Dobry & Gazetas (1986), Pais & Kausel (1988) and Wolf (1988). The effect of foundation embedment on stiffnesses has been studied numerically, but expressed in a similar way as for the bearing capacity: using closed-form depth-of-embedment

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‘correction’ factors to enhance the stiffnesses over those of the surface footing with the same base geometry (Kausel & Roesset, 1975; Tassoulas & Kausel, 1983; Wolf, 1985, 1988; Gazetas, 1991). The historical developments in this soil–foundation interaction field (albeit from a dynamic point of view) have been presented recently by Kausel (2010), while Doby (2014) has compiled a wealth of solutions obtained with simple, physically inspired methods.

In the last 10 or 15 years, studies of dynamic soil–foundation–structure interaction have explored the idea that under seismic (base) excitation mobilising the ultimate capacity of foundations may be beneficial for the whole system. The concept of ‘rocking isolation’ has thereby emerged as an alternative to the conservative – yet not safer – design of foundations against strong seismic shaking (Pecker, 1998; Martin & Lam, 2000; Faccioli et al., 2001; Paolucci et al., 2008; Gajan & Kutter, 2008; Anastasopoulos et al., 2010a, 2012; Gazetas, 2015). As a result of these studies, the response of foundations under cyclic loading mobilising the ultimate capacities has become a key part of the seismic foundation design – bridging the gap between static and dynamic analysis.

Studies of the ultimate response of embedded foundations refer mostly to strip and circular foundations, the external surfaces of which are perfectly ‘glued’ to the soil (Bransby & Randolph, 1999; Poulos et al., 2001; Yun & Bransby, 2007; Gourvenec, 2008). These studies have focused mainly on the undrained bearing capacity under combined loading, modelling the soil as an elastoplastic continuum. Recent studies treating the soil as an inelastic and tensionless Winkler medium include those by Allotey & Naggar (2003) and Gerolymos & Gazetas (2006a, 2006b, 2006c).

The aim of the current study is to extend previous work in an effort to quantify the inelastic response of embedded foundations. A series of three-dimensional (3D) finite-element (FE) analyses are reported treating the soil as an elastoplastic continuum in order to investigate

(a) the effect of embedment on the undrained capacity of square-in-plan foundations under combined loading for a fully bonded contact (FBC) of the foundation with the soil; simple analytical solutions offer enhanced understanding of the mechanics of the problem
(b) the effect of having a tensionless interface of limited shear capacity (TSI) on the failure envelopes of an embedded foundation
(c) the effect of embedment on the elastic stiffnesses in rotation and horizontal translation of the foundation in FBC with the soil; and the relationship of the effects of embedment on (i) bearing capacity and (ii) stiffness
(d) the rotational (‘rocking’) response of simple but slender structures on embedded foundations under monotonic and cyclic displacement-controlled loading, driving them to large deformations after their moment capacity is reached.

Square foundations of various embedment ratios \(0 \leq D/B \leq 1\) are considered under undrained loading conditions. As depicted in Fig. 1 the problem is analysed in two different phases. First (Fig. 1(a)), the bearing capacity of the soil–foundation system is examined under combined monotonic MQ loading applied directly at the foundation base level, not accounting for second-order phenomena. Failure envelopes are predicted in typical sections of the 3D load space.

In the second part (Fig. 1(b)), the entire soil–foundation–structure system is modelled, where the superstructure is represented by a rigid beam element with a lumped mass at its top. The system is subjected to monotonic and cyclic horizontal loading applied at the level of the lumped mass. The load combinations are compatible with the slenderness of the superstructure (i.e. the beam length) and \(P-\Delta\) effects are taken into account. Emphasis is placed on the failure mechanisms, the comparison between the effects of embedment on stiffness and capacity, the cyclic behaviour and the residual deformations of the foundation.

Finite-element modelling

A series of 3D FE analyses are performed using Abaqus (2008). A typical 3D FE mesh, taking advantage of problem symmetry, is shown in Fig. 2. Soil and foundation are modelled with eight-noded hexahedral brick-type elements,
non-linear for the former and elastic for the latter. Fully integrated first-order isoparametric elements are employed for this purpose, in which the volume changes at Gauss points are replaced by the average volume change of the element. Such selective reduced integration (because the order of integration is reduced in selected terms) helps prevent mesh locking, providing accurate solutions even in (nearly) incompressible materials. Its efficiency has been verified in previous publications (e.g. Gazetas et al., 2013), and through initial sensitivity and benchmark analyses for the specific problem analysed herein.

The soil is a homogeneous clay stratum modelled as an elastoplastic material (with parametrically variable undrained shear strength $S_u = 75–150$ kPa, $E/S_u = 1800$, $γ = 20$ kN/m$^3$) while the foundation is modelled invariably as a rigid body. The base of the model is fixed in all three coordinate directions. Special interface elements are used for the soil–foundation interface, modelling two different conditions as ‘(artistically) sketched in Fig. 3.

(a) Fully bonded contact (FBC): this is where the foundation is assumed to remain in perfect contact with the surrounding soil (Fig. 3(a)). Such contact has infinite tensile and shear capacities, preventing separation and slippage of the foundation from the soil.

(b) Tensionless sliding interface (TSI): in this case, separation (gapping) of the foundation from the soil, as well as slippage at the soil–foundation interface, are permitted (Fig. 3(b)). The latter obeys in total stress analysis Coulomb's friction law with $\phi_\alpha = 0$ and $c_\alpha = α S_u$, while detachment and uplifting arise from the tensionless interface behaviour. In the analyses presented herein, the adhesion coefficient, $α$, was set equal to 0.80.

**Soil modelling**

The non-linear soil behaviour is modelled through a simple kinematic hardening model with Von Mises failure criterion and associated flow rule. The evolution law of the model consists of two components: a non-linear kinematic hardening component which describes the translation of the yield surface in the stress space (defined through a 'back-stress' parameter $β$), and an isotropic hardening component which defines the size of the yield surface $σ_0$ as a function of plastic deformation. The pressure-independent yield surface of the model according to the Von Mises failure criterion is defined through the following function $F$

$$F = f(σ – β) – σ_0 \tag{1}$$

The model hinges on three parameters: the elastic Young’s modulus $E$, the ultimate strength $σ_u$, and the yield stress $σ_0$. The evolution of the kinematic component of the yield stress is defined as

$$\dot{σ} = C \frac{1}{σ_0}(σ – α){ε}^p – μσ^{pl} \tag{2}$$

where $C$ is the initial kinematic hardening modulus ($C = σ/ε_0 = E(1 + ν)G_0$) and $γ$ is a parameter determining the rate of decrease of the kinematic hardening with increasing plastic deformation. In the case of clay, the maximum yield stress can be defined as

$$σ_y = \sqrt{3}S_u \tag{3}$$

And since $σ_y = C/γ + σ_0$, parameter $γ$ can be expressed as

$$γ = \frac{C}{\sqrt{3}S_u – σ_0} \tag{4}$$

Parameter $σ_0$ controls the initiation of non-linear behaviour and is defined as a fraction $λ$ of the yield stress $σ_y$

$$σ_0 = λσ_y \tag{5}$$

Finally, $C$ corresponds to the Young’s modulus for very small strains.

Details on the calibration of the model parameters from published experimental $G – γ$ and $ζ – γ$ curves (Vucetic & Dobry, 1991), and the reasonably satisfactory comparisons of its predictions against centrifuge model tests and analytical solutions, can be found in papers by Anastasopoulos et al. (2011, 2012) and Adamidis et al. (2014). In all the analyses, the calibration of parameters was based on the Vucetic & Dobry (1991) $G – γ$ curves for plasticity index of the soil, $PI = 30$.

**Loading conditions**

In the first phase (see Fig. 1(a)), displacement-controlled monotonic loading was applied directly at the centre of the foundation base. A fraction of the ultimate vertical bearing capacity $N_b$ was imposed as a direct vertical force $N_b$ followed by a displacement probe of constant ratio $μ/θ$ until the foundation load does not vary with increased displacement. In the words of Yun & Bransby (2007): ‘once the failure envelope is reached, each loading path travels around the failure envelope until it reaches a termination point where the direction of the tangent to the failure envelope matches the prescribed displacement ratio’.

As associative flow rule was assumed in the constitutive model, the failure envelope serves also as the plastic potential surface, with the direction of the incremental plastic displacement vector at failure being perpendicular to that surface.

It is noted here that the selection of the load reference point affects significantly the shape of the failure envelope produced. In this paper the failure envelopes refer to the centre of the foundation base, in order to be compatible with numerous already published results.

**Model validation**

The constitutive model has been validated against physical model tests (centrifuge, 1g, and large scale) for: (a) surface and embedded foundations subjected to cyclic loading and seismic shaking (Anastasopoulos et al., 2011, 2012); (b) piles subjected to cyclic loading (Giannakos et al., 2012); (c) bar-mat retaining walls subjected to seismic shaking (Anastasopoulos et al., 2010b); and (d) circular tunnels subjected to seismic shaking (Bilotta et al., 2014; Tsinidis et al., 2014). The latter two offer an independent validation, as part of a carefully planned round robin numerical test using centrifuge model tests conducted at the University of Cambridge (Lanzano et al., 2012). Despite its simplicity, the model was shown to perform very well, in many cases outperforming more sophisticated models (Bilotta et al., 2014).

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**Fig. 3. Sketched representation for the meaning of soil–foundation interface conditions: (a) FBC; (b) tensionless and potentially sliding interface – TSI**
The model is further validated herein against published failure envelopes for embedded foundations subjected to combined $NQM$ loading. As mentioned earlier, most of the available solutions are limited to plane-strain conditions. For a meaningful comparison, the 3D FE model of Fig. 2 is transformed into an equivalent two-dimensional (2D) model by considering only an appropriate ‘slice’. A FBC is considered: sliding or detachment is disallowed. Several analyses were conducted using different $u/\theta$ ratios to produce the failure envelope for a certain plane in $NQM$ loading space. Fig. 4(a) shows the load paths for different displacement probes in comparison with the respective failure envelope of Gourvenec (2008) for an embedment ratio $D/B = 1$. To the present authors’ satisfaction, the terminal points of the load paths fall on top of, or sufficiently close to, Gourvenec’s envelope. An additional validation for $D/B = 0.5$ against the failure envelope of Vulpe et al. (2014) for circular foundations and Gourvenec (2008) for strip foundations is shown in Fig. 4(b).

NUMERICAL AND (SOME) ANALYTICAL RESULTS

The undrained bearing capacity of rigid embedded strip foundations under combined $NQM$ loading has been extensively investigated by, among others, Yun & Bransby (2007) and Gourvenec (2008), who considered plane-strain (2D) conditions and a fully bonded soil–foundation contact. The current study extends these results: (a) to the 3D problem of a square embedded foundation; and (b) to the case of a tensionless interface which accommodates both separation and sliding (the latter when a certain fraction, $\alpha = 0.80$, of the soil undrained shear strength is reached: $f_s = 0.80 S_u$); (c) to the comparative study of the influence of embedment on stiffness against that on ultimate capacity; and (d) to the response under quasi-static cyclic loading coming from lateral loading on top of a slender structure.

FULLY BONDED CONTACT

Lateral bearing capacity and elastic stiffness

Numerical results from the 3D FE analysis (FEA) of square-in-plan embedded foundations are presented and compared with those of the corresponding plane-strain analyses. Thanks to the additional two side surfaces, parallel to the direction of loading, the square foundation enjoys increased soil resistance beyond the one experienced by an equivalent embedded strip (2D) foundation for all modes of loading: vertical, horizontal and moment. Furthermore, the
unavoidable three-dimensionality of the failure mechanism elicits increased soil resistance, as explained below.

**Horizontal translation (without rotation).** A simple failure mechanism aimed merely at illustrating the origin of the effect of embedment is depicted in Fig. 5. The lateral bearing capacity $Q_{\text{max}}$ of a square embedded footing is examined under pure horizontal translation (rotation $\theta = 0$). In two dimensions the foundation ultimate horizontal capacity $2dQ_{\text{max}}$ could be obtained roughly as the superposition of the ultimate base shear ($2dQ_{\text{base}}$), which is essentially equal to the ultimate shear force $2dQ_{\text{uo}}$ of a surface footing ($2dQ_{\text{uo}} \approx 2dQ_{\text{base}} \approx S_u A_{\text{base}}$) and the passive-minus-active force ($2dP_{\text{passive}} - 2dP_{\text{active}}$) obtained from Coulomb-type plane-strain analysis (here assuming a circular sliding surfaces motivated by the FE observations of, say, Fig. 9). For such type of analysis (but using a rigorous method) Gourvenec (2008) has fitted the following expression to the numerical results for the maximum horizontal load as a multiple of the corresponding load of the surface footing

$$2dQ_{\text{max}} \approx 2dQ_{\text{uo}}[1 + 5.56(D/B)]$$

Interestingly, for the elastic horizontal stiffness of a strip foundation, the effect of embedment is much less significant (Gazetas, 1983, as modified recently by Lekkakis, 2012)

$$2dK_{\text{H}} \approx 2dK_{\text{HO}}[1 + 1.2(D/B)]$$

For example, for $D/B = 1$, the increase in stiffness due to embedment is only about 1/3 of the corresponding increase in bearing capacity. Here are two possible culprits of such a greater effect of embedment on ultimate load as opposed to its effect on stiffness.

(a) The most significant: the sidewall (vertical) surfaces at small elastic deformations invoke a disproportionately small soil reaction; because each sidewall can be viewed as nothing more than a vertical strip at the edge of a homogeneous quarter-space (i.e. the space between the horizontal surface from the footing edge to infinity along the $x$-axis, and the vertical axis passing along the sidewall, as sketched in the insert of Fig. 5); that strip is loaded horizontally (i.e. perpendicularly to its surface). Recall that on a homogeneous half-space the normal

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**Fig. 5.** (a) Estimation of maximum horizontal capacity of three-dimensional embedded foundations with finite plan dimensions: isolation of a slice of soil which slides against the surrounding soil. (b) Sketch illustrating the role of the sidewall that is normal to the direction of loading.
stiffness, $K_0$, of a strip is vanishingly small. It will be even smaller on a quarter-space. This is because on a homogeneous half-space an infinitely long and uniformly loaded strip leads exactly to infinite displacement – vertical or horizontal, depending on the loading direction (normal or shear). Hence $K_0$ (or $K_{1H}$) = 0. This is an unequivocal result (a fact) on the theory of elasticity. Naturally, in real life neither a homogeneous half-space (i.e. a homogeneous medium of infinite depth) nor a strip loading (i.e. of infinite extent) can exist. The soil will become inhomogeneous or meet a stiff soil layer or rock at some depth, $H$. In such a case $K_0$ decreases as $H/B$ increases. But this has nothing to do with the idealisation of the embedded foundation – the side soil extends to infinity in the horizontal direction. By contrast, the (Coulomb-type) sliding wedge is of limited extent, and the resulting passive resistance is relatively huge (requiring of course a very large deformation). The ‘true’ failure mechanism from the FEA, to be discussed later in this paper, elucidates this point. Hence the much greater contribution of the side walls to the overall foundation capacity rather than stiffness.

(b) The small-deformation elastic regions below the base and in front of the sidewalls affect each other (a base-to-side interaction), thus reducing the overall stiffness to a smaller value than what the summation of their effects would provide; at failure, however, the mechanism comprises three more-or-less independent sliding surfaces, and therefore the ultimate load is essentially the sum of the respective three resistances.

Now, under 3D conditions, two additional foundation sidewalls are available, parallel to the loading plane resisting in pure shear the horizontal translation. The maximum additional force is given by

$$T_{	ext{sidewall}} \approx S_u A_{	ext{sidewall}} = 2 \times B D S_u$$

(8)

Furthermore, the simple mechanism of Fig. 5 consists of two quarter-cylinder Coulomb-type sliding wedges, as would have also been appropriate for the 2D problem; but now, additionally, the four sides of the two cylinders experience shear stresses $S_u$ from the outside un-deforming soil. Hence the passive-minus-active force increases by

$$F_{	ext{sidesoil}} \approx S_u \pi D^2/4 \times 4 = S_u \pi D^2$$

(9)

Therefore, assuming for simplicity that the three partial mechanisms act independently, an upper-bound estimate for the maximum horizontal force of the square-in-plan embedded foundation would be

$$Q_{\text{max}}/Q_{\text{ult}} \approx Q_{\text{max}}/B^2 S_u$$

$$\approx \left[ 1 + 5.56((D/B)) \right] + \left[ 2(D/B) \right] + \left[ \pi(D/B) \right]^2$$

(10)

in which $Q_{\text{ult}} \approx B^2 S_u$ is the ultimate horizontal capacity of a surface square foundation. Fig. 6 compares the $Q_{\text{max}}/B^2 S_u$ plotted against $D/B$ relationship from the FEA with the above approximate expression. Given the simplicity of the failure mechanism (hardly optimised) for the passive and active resistance, equation (10) is surprisingly (and perhaps fortuitously) accurate.

Admittedly, however, whereas equation (10) is independent of the vertical load $N$ acting on the foundation, the FEA gives results that are slightly affected by the vertical load $N$. This explains the approximate equality sign between

$$Q_{\text{max}} \approx 11.5$$

which is a very substantial effect, indeed. Again, compare the above with the elastic horizontal stiffness of a square-in-plan foundation with the same embedment $D/B = 1$ (Gazetas, 1991; Lekkakis, 2012)

$$K_{1H}/S_u \approx \left[ 1 + 0.21((D/B)0.5) \right] \left[ 1 + 1.57((D/B)0.8) \right]$$

$$\approx 3.1$$

(13)

in which $K_{1H}$ is the elastic horizontal stiffness of the square surface foundation. Notice that, once more, the effect of embedment on the horizontal stiffness is an even smaller fraction (about 1/4) of the effect on ultimate load. The explanation is as above for the 2D problem, but with four rather than two sidewall surfaces participating now. Fig. 7(a) illustrates the (purely horizontal) failure mechanism for two embedment ratios (0.2 and 1), in the form of the vectors of displacement, with superimposed as dark shadows the regions of large plastic strains. It is interesting to observe that the plastic strains tend to localise in well-defined shear bands. In such cases, the problem can be mesh dependent and a sensitivity analysis is required. Such a sensitivity study was performed as part of the present study, leading to the selection of the adopted FE mesh.

Rotation about the centre of the base. The overturning moment, $M_{\text{max}}$, producing a rotational failure without horizontal displacement, is similarly affected by the presence of the sidewalls that are perfectly bonded to the surrounding soil. Fig. 8 plots the variation of $M_{\text{max}}/BAS_u$ against $D/B$ from the FEA.

To compare the 3D results of this figure with the 2D effects of embedment, recall again the plane-strain analysis of Gourvenec (2008), who has fitted the following expression to her numerical results for the maximum
moment

\[ 2dM_{\text{max}} \approx 2dM_{\text{uo}} \left[ 1 + 1.59(D/B) + 3.22(D/B)^2 \right] \]  
(15)
where \( 2dM_{\text{uo}} \) is the ultimate moment of the fully bonded surface strip foundation.

As an example, for \( D/B = 1 \), equation (15) yields an ultimate moment ratio \( \frac{2dM_{\text{max}}}{2dM_{\text{uo}}} \approx 5.8 \). This can be compared with the effect of embedment on the elastic rotational stiffness of a strip foundation

\[ \frac{2dK_{R}}{2dK_{RO}} \approx 1 + 2(D/B) = 3 \]  
(16)
where the rotational stiffness of the surface strip (per unit running length) is

\[ 2dK_{RO} = \pi GB^2/[8(1 - \nu)] \]  
(17)

The effect of embedment on stiffness is now more substantial, but still about 1/2 of the effect on ultimate rotational capacity. The cause of this increased relative importance of embedment on stiffness will be discussed and explained below.

The 3D square-in-plan embedded foundation studied here would potentially benefit from the following additional contributions to resistance over and above the embedded strip footing.

(a) The base now is square rather than strip; the failure mechanism is different and the ultimate pure moment of the fully bonded footing is approximately

\[ M_{\text{uo}} \approx S_uBA \]  
(18)

(b) The two sidewalls that are parallel to the direction of loading undergo a torsional-type movement about the horizontal axis passing from the centre of the base, thus contributing by an amount

\[ M_{T,\text{sidewall}} \approx S_uBD \times D/2 \times 2 = S_uBD^2 \]  
(19)

(c) The two sidewalls that are perpendicular to the loading direction undergo a movement which can be decomposed into a vertical and a horizontal translation of their centre, equal to \( \theta B/2 \) and \( \theta D/2 \), respectively, and rotation \( \theta \) around their centroidal horizontal axis (rocking). Of these, the vertical movement invokes shear resistance and its ultimate value \( A_{\text{sidewall}}S_u \) is the same, equal to \( 2 \times BDS_u \), for the strip and the square. Hence no correction above what applies to the strip is needed. For the rotation and the horizontal compression, the aforementioned correction shape factor of 1.2 would arguably be an upper bound above what applies to the strip. Hence, assuming for now independence of each contributing movement, gives as an average the factor: \( (1 + 2 + 1.2)/3 \approx 1.13 \). In fact, since some interplay among these three components of ultimate resistance is unavoidable, this small correction factor will be ignored altogether, accepting that the second factor on the right-hand side of

\[ \frac{M_{\text{uo}}}{M_{\text{uo}} + M_{\text{T,side}}} \approx \frac{5.8}{5.8 + 3.8} \approx 0.63 \]
Gourvenec’s equation (15) adequately describes the contribution of the two perpendicular sidewalls. Hence, as a crude approximation, the maximum moment is

\[ M_{\text{max}} \approx M_{\text{uo}} [1 + 1.59(D/B) + 4.22(D/B)^2] \]

which for the example of \( D/B = 1 \) gives for the normalised moment: \( M_{\text{max}}/M_{\text{uo}} \approx 6.8 \). In Fig. 8 the FEA gives a larger value, 7.5. Given the crudeness of the above simple upper-bound analysis, its performance is quite acceptable.

Again it is worth comparing the above with the effect of embedment on the elastic rotational stiffness of the foundation for \( D/B = 1 \). From the aforementioned publications

\[ K_R/K_{RO} \approx \left[ 1 + 1.3(D/B)/(1 + (D/B)) \right] \approx 3.6 \]

It is quite interesting to notice that the effect of embedment on stiffness for this mode of loading is about 0.60 times the corresponding effect on the (‘true’) bearing capacity.

Table 1 compiles the above comparisons between the aforementioned effects of embedment on capacity and stiffness for horizontal and rotational loading. It is clear that in both 2D (strip) and 3D (square) situations

(a) the rotational stiffness increases approximately by 1/2 to 2/3 as much as the increase in maximum (overturning) moment
(b) the horizontal stiffness increases approximately by only 1/4 to 1/3 as much as the increase in maximum (horizontal) force.

A simple explanation for the relatively increased effect on the rotational stiffness could be unravelled from the previous discussion: any rectangular footing of width \( B \) on a halfspace subjected to moment loading (in rocking or torsion) affects only a limited volume of soil; in the traditional geotechnical jargon, it produces elastic ‘stress bulbs’ of very limited

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extent – one-fourth (1/4) to one-eighth (1/8) of the width $B$ (Poulos & Davis, 1974). In contrast, force loading (in shear or normal mode) produces ‘stress bulbs’ that go much deeper – one to two times the width, respectively. But smaller affected area means smaller total displacements–rotations, and therefore larger stiffnesses. Hence, the sidewalls in rotational modes exhibit much greater stiffness than in translational modes; this explains the trend of Table 1. The difference between the two modes of loading is even larger in the 2D plane-strain problem. The reason: the sidewalls are (infinitely long) strips in this case, possessing negligible stiffness in force loading (theoretically, zero), but offer appreciable stiffness in rocking about the long horizontal axis.

### Table 1. Comparison of the effects of embedment on ultimate loads and elastic stiffnesses for square ($Q, K$) and strip ($2dQ, 2dK$) foundations: results of FEA. $D/B = 1$

<table>
<thead>
<tr>
<th>$2dQ_{\text{max}} / 2dQ_{\text{uo}}$</th>
<th>$2dK_{H}/2dK_{HO}$</th>
<th>$Q_{\text{max}} / Q_{\text{uo}}$</th>
<th>$K_{H}/K_{HO}$</th>
<th>$2dM_{\text{max}} / 2dM_{\text{uo}}$</th>
<th>$2dK_{R}/2dK_{RO}$</th>
<th>$M_{\text{max}} / M_{\text{uo}}$</th>
<th>$K_{R}/K_{RO}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>6.6</td>
<td>2.3</td>
<td>11.5</td>
<td>3.1</td>
<td>5.8</td>
<td>3.0</td>
<td>7.5</td>
<td>3.6</td>
</tr>
</tbody>
</table>

Fig. 10. Graphical representation of $MQ$ interaction envelope with the direction of plastic deformations at four key points. Four offset diagrams show displacement vectors and concentration of shear strains revealing the failure mechanisms at these points. FBC, $D/B = 0.2$. The solid dots in the pictures represent the instantaneous rotation pole of the foundation. (The pole is at infinity for the horizontal translation case, point C, and ‘almost’ infinity at point D.)

### Failure Envelopes

**MQ interaction (with $N = 0$)**

So far, aspects of the maximum horizontal load $Q_{\text{max}}$ and of the maximum overturning moment $M_{\text{max}}$ have been studied separately. Of course, it was a combination of $Q$ and $M$ that led to these two limit values: a moment was necessary to nullify the rotation as required for $Q_{\text{max}}$ and a horizontal force to nullify the horizontal displacement as required for $M_{\text{max}}$. But these were merely two $MQ$ combinations out of infinite possibilities. The full interplay between the limiting values of $Q$ and $M$, comprising all combinations, is portrayed in Figs 9(a) and 10, in the normalised form: $M/BAS_u$ plotted against $Q/AS_u$ for two values of $D/B$, 1 and 0.2. The vertical load, $N$, which also affects the $QM$ interaction, is taken as 0 in both figures. Figure 9(b) is only meant to illustrate the fact that the shape of the failure envelopes depends strongly on the point which $M$, $Q$, $u$ and $\theta$ refer to. Of course, the information conveyed by each of the three different sets of envelopes (interaction diagrams) for the three reference points (base, middle, top) is exactly the same.

In Figs 9(a) and 10 the location of $Q_{\text{max}}$ and $M_{\text{max}}$ is shown at the two extremes of the envelope, where the normal lines to the envelope (the ‘failure surface’) are parallel to the horizontal and vertical axes, respectively. This is a consequence of the associated flow rule adopted in the present soil constitutive model (preservation of normality). Also indicated in the figures are the two limit values of the purely horizontal-force ($M = 0$) and purely overturning-moment ($Q = 0$) loading, denoted as $Q_{\text{ult}}$ and $M_{\text{ult}}$, respectively (adopting the terminology of Gourvenec). The normal vectors to the envelope at these latter locations point to the
negative $M$ and $Q$ axes, respectively – indicative of negative (counter-clockwise) rotation and negative (on the $x$-axis) displacement, respectively. This behaviour can be visualised in the inserted four snapshots of the vectors of displacement (with the concentration of plastic shear strains superimposed as shadows), for each of the four limit loads discussed here. They correspond to pure horizontal load, pure horizontal translation, pure moment and pure rotation, respectively.

Figures 9(a) and 10 reveal a significant effect of embedment, not only on the values of the various limit loads but also on the shape of the envelope; the latter becomes increasingly skewed with increasing (relative) depth, $D/B$. As a result, the maximum moment capacity $M_{\text{max}}$ (which occurs in the presence of positive horizontal load, applied always at the base) increases disproportionately more than the increase of $M_{\text{ult}}$. Indeed, the ratio $M_{\max}/M_{\text{ult}}$ attains (approximately) the values of

- (a) 1.75 for $D/B = 1$
- (b) 1.10 for $D/B = 0.2$

Similar are the trends of the $Q_{\max}/Q_{\text{ult}}$ ratio, which is equal to

- (a) 1.60 for $D/B = 1$
- (b) 1.02 for $D/B = 0.2$

For the surface foundation, that is, $D/B = 0$, both ratios approach 1, as there is negligible coupling of $M$ and $Q$ under fully bonded contact and $N = 0$, as considered here.

To further elucidate the role of embedment on the shape of the $QM$ failure envelopes, Fig. 11 compares the four envelopes (for $D/B = 0, 0.2, 0.5$ and 1) as plotted in the normalised coordinate system: $M/M_{\text{ult}}$ plotted against $Q/Q_{\text{ult}}$.

It is also worth observing the failure mechanisms at $Q_{\text{ult}}$ and $M_{\text{ult}}$ in the inserts in Figs 9 and 10. Under pure moment loading, at $M_{\text{ult}}$, applied at the foundation base level, a scoop mechanism is observed with its rotation pole (indicated with the solid dot in the figure) located approximately $D/2$ above the base, that is, at the foundation centre of gravity. The foundation response is therefore mainly rotational, but it is also accompanied by a negative horizontal translation. The latter, as indirectly evidenced by the oblique intersection of the failure envelope with the moment ordinate axis, is substantial only for the deeper foundation ($D/B = 1$); it is of marginal importance for the shallow foundation ($D/B = 0.2$), and would hardly exist for a surface foundation.

On the contrary, under purely horizontal loading, at $Q_{\text{ult}}$, a reverse scoop mechanism is formed with its centre of rotation moving up near the soil surface – nearly a pendulum. Thus, failure consists of (nearly equally important) horizontal translation and counter-clockwise rotation. In this case, the obliquity of the intersection of the envelope with the $Q$ axis is appreciable even for the shallow foundation ($D/B = 0.2$).

At the limiting value $Q_{\text{ult}}$ of exclusively horizontal displacement, the failure mechanism is quite different: sliding of the foundation–soil system develops, with active and passive failure on the front and back side of the foundation, and shear at the base and at the two parallel sidewalls, as already discussed. (Notice in passing the nearly circular (in cross-section) failure mechanisms on both the active and passive sides, justifying the simplified mechanism invoked in the preceding section to obtain equation (9).) At the limit value for pure rotation, $M_{\text{ult}}$, the failure mechanism is composite: spherical failure surface below the base, active and passive wedges on the upper part of the normal sidewalls, and torsional shear on the parallel sidewalls.

QN interaction (with $M = 0$)

The significance of the vertical load for the ultimate value $Q_{\text{ult}}$ of a solely horizontally lateral load ($M = 0$) is portrayed in Fig. 12, for four values of $D/B$: 0, 0.2, 0.5, 1. More specifically: in Fig. 12(a) in the form of $Q_{\text{ult}}/AS_{u}$ plotted against $N/AS_{u}$ and in Fig. 12(b) in the normalised form $Q_{\text{ult}}/Q_{\text{ult},N=0}$ plotted against $N/N_{\text{ult},Q=0}$. Also plotted in Fig. 12(a), for a mere comparison, are the corresponding 2D curves of Gourvenec (2008) for the strip foundation (in which case $A = B$). Observe in Fig. 12(b) that the differences between the normalised curves as functions of $D/B$ are barely if at all distinguishable; also indistinguishable would be the square (3D) and strip (2D) normalised curves (and which are thereby not shown) – proofs of nearly identical shapes of envelopes for both types of foundations (strip and square) and all embedment ratios!

Note in Fig. 12 that for low vertical loads compared to the vertical capacity, $N/N_{\text{ ult}, Q=0} < 0.5$, or equivalently for safety factors $FS_{V}$ against vertical bearing capacity mobilisation exceeding 2 (a most frequent situation in practice), the horizontal capacity remains almost constant, regardless of the axial force magnitude. This is true for all examined embedment ratios. However, an abrupt reduction in load-carrying capacity is noticed at higher vertical loads, with $Q_{\text{ult}}/Q_{\text{ult}, N=0}$ dropping to about 0.5 for an (admittedly very small) $FS_{V} = 1.11$ (i.e. $N/N_{\text{ult},Q=0} = 0.9$).

MN interaction (with $Q = 0$)

The significance of the vertical load for the ultimate value $M_{\text{ult}}$ of a purely moment loading ($Q = 0$) is portrayed in Fig. 13, for four values of $D/B$: 0, 0.2, 0.5, 1. More specifically: in Fig. 13(a) in the form of $M_{\text{ult}}/ABS_{u}$ plotted against $N/AS_{u}$ and in Fig. 13(b) in the normalised form $M_{\text{ult}}/M_{\text{ult},N=0}$ plotted against $N/N_{\text{ult},M=0}$. Also notice in Fig. 13(a), for a mere comparison, the corresponding 2D curves of Gourvenec (2008) for the strip foundation (in which case $A = B$). Again, in Fig. 13(b) the strip (2D) normalised curves would be indistinguishable from the 3D ones if they were shown – proof of nearly identical shapes of the envelopes for both types of foundations and for all studied embedment ratios!

Now the vertical load plays a slightly greater role than for the horizontal capacity. Again, at low values of $N/N_{\text{ult}, M=0}$ the effect of $N$ on the $M_{\text{ult}}$ is negligible, but the limit now is $N/N_{\text{ult}, M=0} < 0.3$ rather than 0.5. This is true also for all examined embedment ratios. However, an abrupt reduction
in load-carrying capacity is noticed at higher vertical loads, with $M_{ult}/M_{ult,N-0}$ dropping to about 0·5 for the very small $FS_v=1·25$ (i.e. $N/N_{ult,M-0}=0·8$).

TENSIONLESS, POTENTIALLY SLIDING INTERFACE (TSI)

MQ envelopes with $N \approx 0$

A very substantial reduction of the effects of embedment ensues when the five interfaces of the foundation with the soil are incapable of transmitting (net) tensional normal stresses and shear stresses that exceed the adhesive stress $f_s=\alpha S_u$. Hence, separation as well as sliding of parts of the walls or the base from the surrounding and underlying soil is a possibility.

Figure 14 shows the $MQ$ envelope for the $D/B=1$ for a foundation with TSI, along with the snapshots of the displacement vectors and the high plastic shear strain ‘shadows’. This figure should be compared with Fig. 9 for the same foundation but with FBC, to get an idea of the effects of separation and sliding. Notice the following points.

(a) For all combinations of $Q$ and $M$ the limiting values of Fig. 14 decrease due to TSI, and the skewedness of the envelopes seen in Fig. 9 substantially diminishes, being barely distinguishable for the smallest embedment ratios, $D/B=0·2$, examined. (b) The failure mechanisms at the four limit points (A–D) seem to consist of only a part of the mechanisms of the FBC analysis of Fig. 9 since, as a consequence of separation, failure is restricted locally in the close vicinity of the particular interface. For instance, at point C of $Q_{max}$ all the back side of the wall loses contact with the soil, and thus no active wedge develops.

At point B of the $D/B=1$ foundation, corresponding to $M_{max}$, the TSI mechanism in front and back of the normal-to-loading-direction sidewalls, as well as the one below the base, are almost half (in a qualitative sense) of the FBC mechanisms. The separation of the back wall is due to the combined effects of the clockwise moment $M=M_{max}$ and the positive $Q$ required for keeping the base horizontal translation equal to 0.

On the other hand, at the $M_{max}$ point B of the $D/B=0·2$ foundation, the TSI mechanism is altogether different: instead of the scoop mechanism of the FBC foundation of Fig. 9, a Brinch–Hansen-type (1953) wedge is evident.

Figure 15 compares the dimensionless $MQ$ envelopes from the FBC and the TSI analyses, for the four $D/B$ ratios (0, 0·2, 0·5, 1). This graph refers to a relatively small but non-zero value of vertical load ($N/N_{ult}=0·25$) to allow comparison with the $D/B=0$ foundation. In the case of the TSI analysis,
such comparison would not be possible for $N/N_{ult} = 0$. This comparison shows that the TSI curves are not only much ‘contracted’ compared to the FBC curves, but their shapes are also quite different, with the skewness having almost disappeared.

Moment against angle of rotation; soil reactions ($u=0$, $N \neq 0$)

Although the ultimate capacity (moment or horizontal) of an embedded foundation under undrained conditions is only marginally affected by the vertical load, $N$, as long as the latter does not exceed 1/2 to 1/3 of the vertical bearing capacity, $N_{ult} = N_{ult,M=Q=0}$, it is worth studying the complete moment–rotation curves for a very large factor of safety and for more usual factors of safety: $FS = N_{ult,M=Q=0}/N = 14$ and 3, representing lightly loaded foundations (or strong soil) and heavily loaded foundations (or weak soil), respectively. The two curves are plotted in Fig. 16. In both cases, the total resisting moment initially increases linearly with rotation angle. The corresponding zero-angle-of-rotation stiffnesses are only very slightly different (difficult to discern in the scale of the figure), because, as was quantitatively shown by Gazetas et al. (2013) and Adamidis et al. (2014), the heavily loaded (in the vertical direction) foundation induces more plastic deformation in the underlying soil and hence depresses slightly even its lateral effective stiffness. But as soon as initiation of either separation/sliding at the interfaces, or soil yielding, or both, take place, a gradually softening rocking behaviour is observed until the resisting moment reaches its maximum value. Overall, the behaviour of the two foundations is quite similar.

The distribution of soil reactions on the foundation walls is also portrayed in this figure, for four different stages of loading. Only normal stresses are shown on the two vertical walls (front and back) and the base. At loading points A, compressive and tensile stresses develop along the foundation periphery. Superimposed on the static (gravity-induced) compressive earth pressures they result in an increase of normal stresses in the compression area and a decrease of stresses in the tension area. In the upper half of the back wall, because of the interface’s inability to sustain pure tension, separation seems to have already occurred at very small angles of rotation. By the time points B are reached, the above picture has become clearly distinguishable.

Fig. 14. $MQ$ interaction (for $N/N_{ult}=0$): failure mechanisms at key points of the failure envelope of a $D/B=1$ foundation having a tensionless sliding interface (TSI) with the soil. The solid dots represent approximately the instantaneous rotation pole, which for purely horizontal translation (point C) is at infinity.

Fig. 15. Effect of interface conditions: failure envelopes in $MQ$ load plane for $N/N_{ult}=0.25$: (a) FBC, compared to (b) tensionless sliding interface (TSI)
As the load increases (points C), ultimate soil reactions seem to have been mobilised near the top of the front sidewall and near the bottom of the back sidewall, similarly on the two foundations. The separation of the back sidewall from the soil has gone deeper in both cases. However, the response at the bases of the heavily and the lightly loaded foundations differ substantially: uplifting of half the base takes place under light (vertical) loading, while compression is still transmitted under the whole heavily loaded base. This is reminiscent of the response of surface foundations that has been amply demonstrated both theoretically and experimentally (Paolucci et al., 2008; Anastasopoulos et al., 2010a, 2012; Deng & Kutter, 2012): uplifting dominates with large FSV, whereas soil yielding and failure mechanisms dominate with small FSV.

With further increase in imposed rotation, the contact areas of the sidewalls become somewhat larger (a small reversal of the previous trend) and the soil reactions increase. The base remains almost in full contact (even if barely) under the heavy load. Also unchanged is the area of uplifting of the base under the light load (half of the base remaining in contact). One can appreciate that there is a more-or-less stable distribution of soil stresses beyond point D, in both cases.

**CYCLIC LOADING AT LARGE DEFORMATIONS**

Rotational response for slender structures: surface foundation

All the results in this section refer only to the non-linear TSI. In general, the response of foundations under strong lateral cyclic loading coming from the superstructure involves rocking, swaying and settlement. Any of these modes can be significant, but for slender structures (say, with $h/B > 1$) rotation dominates over horizontal translation. Thus, to simplify the picture of cyclic response, purely horizontal displacement is not discussed here.

In this section the authors explore the significance of embedment ($D/B$) on the cyclic rotational (‘rocking’) behaviour not just of an isolated foundation, but rather of a foundation supporting a simple rigid structure of weight $W$ located at a height $h$ from the foundation top. With tensionless and potentially sliding soil–foundation interfaces (TSI), in addition to the amplitude of loading, the main parameters to investigate are the vertical static factor of safety, FSV, and the slenderness ratio, $h/B$; recall Fig. 1(b). A rotation-controlled loading is applied to the structure–foundation system. The loading protocol shown in Fig. 17 comprises six cycles of progressively increasing angle of rotation $\theta$ from 0·0025 rad to 0·08 rad (i.e. 2·5 to 80 millirad), covering the range from quasi-linear to highly non-linear response. The rotation is imposed applying a lateral displacement at the mass, as shown in Fig. 1(b). The mass is adjusted appropriately in order to achieve the desired FSV for each case studied. Constant-amplitude multi-cycle tests have also been performed.

![Fig. 16. Illustration of the evolution of stresses at the soil–foundation tensionless sliding interface (TSI) for two different axial loads corresponding to FSV values of 14 and 3: (a) moment–rotation curves with selected four loading levels; (b) stress distribution at these four levels](image)

![Fig. 17. Cyclic loading protocol (for Figs 18–21)](image)
Developed dimensionless moment–rotation \((M/B_{AS})\) and settlement–rotation \((\theta)\) curves now refer to the base of the structure rather than of the foundation, as was the case until now. These curves are presented along with the corresponding dimensionless settlement–rotation \((\theta)\) curves as follows: in Figs 18 and 19 for surface foundations \((D/B = 0)\), and in Figs 20 and 21 for embedded foundations with \(D/B = 0.5\) and 1. Two FSV values are explored: 5 and 2, representing, respectively, lightly loaded structures (or strong soil) and heavily loaded structures (or soft soil). The slenderness ratio, \(h/B\), is varied parametrically (= 1 and 3) for the surface foundations, but is kept constant (= 1) for the embedded ones.

For the surface footings, Fig. 18 for the very slender supported structure \((h/B = 3)\) shows a decline of the monotonic (‘backbone’) curves after the ultimate moment has been reached, for both FSV values. This decline is a consequence of the \(P–\Delta\) effect which stems from the additional moment \(\Delta M = W(h \sin \theta)\) generated by the eccentricity of the weight \(W\) as it displaces horizontally by \(h \sin \theta \approx h \theta\). \(\Delta M\) aggravates the consequences of loading, leading to a post-peak decline of the moment capacity. Of course, the phenomenon is conspicuous only at large angles of rotation; it is barely distinguishable at relatively small angles, even after slight exceedance of the maximum moment, as could be seen in this figure (and will be more obvious in the next figures). And as will be shown later, \(P–\Delta\) is significant mainly for surface (or very shallow) foundations. It is particularly significant for a relatively rigid structure on a rigid base (Makris & Roussos, 2000; Apostolou et al., 2007; Panagiotidou et al., 2012).

A broad similarity is noted in the shape of the two (static) monotonic curves, although the heavily loaded footing (FSV = 2) experiences larger moment capacity, as of course was expected even from the Meyerhof bearing capacity solution, but reaches the overturning angle \((\approx 0.011\) rad, where the curve crosses the 0 axis) sooner than that of the lightly loaded FSV = 5 foundation \((\approx 0.019\) rad at the projection of the curve).

In both cases the monotonic curves clearly ‘envelop’ the loops of the (slow) cyclic tests. Hence, the MQN interaction diagrams presented in the paper are also valid for cyclic loading – an important conclusion. Yet, despite their similarities, the two different FSV foundations behave distinctly differently.

(a) An uplifting-dominated behaviour is noted with the lightly loaded foundation (FSV = 5), as evidenced indirectly from the roughly-S-shaped \(M–\theta\) loops (with negligible residual moment each time the imposed \(\theta\) returns to 0, implying a ‘re-centring’ capability under cyclic excitation), and directly from the upward displacement of the foundation centre seen in the \(w–\theta\) curves. Of course, despite such uplifting (or perhaps because of it) there is a progressive recession (permanent loss of contact) between the soil and the foundation under the edges, which leads to the eventual accumulation of a very small (downward) settlement, and to an appreciable degradation of rotational stiffness. (The latter is seen in the decreased slope of the unloading–reloading lines.)

(b) A sinking-dominated response is noted with the heavily loaded foundation (FSV = 2), as evidenced indirectly from the roughly-O-shaped large inelastic loops in \(M–\theta\), and directly from the accumulating downward displacement in the \(w–\theta\) curves – both phenomena being a consequence of extensive cyclic soil inelastic action.
Fig. 19. Dimensionless moment–rotation and settlement–rotation response of a surface TSI foundation subjected to cyclic lateral loading: (a) $FS_v = 5$; (b) $FS_v = 2$ (grey line corresponds to monotonic backbone curves). $D/B = 0$, $h/B = 1$

Fig. 20. Dimensionless moment–rotation and settlement–rotation response of a $D/B = 0.5$ embedded TSI foundation subjected to cyclic lateral loading: (a) $FS_v = 5$; (b) $FS_v = 2$ (grey line corresponds to monotonic backbone curves). $h/B = 1$
The less slender system with \( h/B = 1 \) is studied in Fig. 19. In general, the response of the two foundations is similar with that of the corresponding foundations of the very slender system in Fig. 18. The main differences are: the \( P-\Delta \) effects are in this case (understandably) less prominent, as \( W(h \sin \theta) \) is reduced; the S-shape of the \( M-\theta \) loops for \( \text{FSV} = 5 \) are not as narrow at the centre; and the accumulated settlements are larger for both foundations (\( \text{FSV} = 2 \) and 5). The latter may come as a surprise (less settlement for the taller structure (of the same weight)), but it is a natural consequence of the type of applied loading; for an imposed large angle of rotation, the total moment atop the foundation is about the same in the two cases (\( h/B = 1 \) and 3); however, the contribution to this moment of the (vertical) weight \( W \) (\( P-\Delta \) effect) is 3 times larger for the taller (\( h/B = 3 \)) structure, and thereby the corresponding moment due to the horizontal force is 3 times smaller. Therefore, the horizontal shear force \( Q = M/h \) developing atop the foundation of the taller system is somewhat smaller than the one coming from the shorter one. Hence in the latter case, the increased \( Q \) leads to greater inelastic action in the soil, resulting in greater residual settlement.

The effect of embedment

Figures 20 and 21 show the effect of embedment (\( D/B = 0.5 \) and 1) for \( h/B = 1 \) (only), and \( \text{FSV} = 5 \) and 2, as before. The following points may be noticed.

(a) The monotonic backbone curves now barely show a sign of \( P-\Delta \) induced degradation up to an angle of rotation of 0.12 rad; just a ‘suspicion’ of a decline with the heavier load.

(b) The shape of the cyclic moment–rotation loops deviates from that of the surface foundations as embedment increases. Lateral loading of the embedded foundation causes excessive plastic deformation of soil on one side with a gap opening on the opposite side. An abrupt decrease in \( M \) upon unloading is attributed to this partial detachment (‘gapping effect’), because momentarily both front and back of the sidewalls are not in contact. Upon reversal, contact is re-established and moment resistance slowly increases. Re-centring of the foundation results in a partial loss of contact between the sidewalls and the soil, as the plastically deformed soil is not restored to its original state.

(c) With increasing number of loading cycles, the soil–foundation gap, alternately forming on either side of the foundation, progressively widens. This results in a decreased resisting moment during unloading and reloading. Notice that the monotonic curve is reached only when the complete passive earth pressures come into action. Thus, an abrupt increase in \( M \) occurs every time the system approaches the angle of rotation corresponding to the amplitude of the previous load cycle.

(d) The differences between the lightly and the heavily loaded system in the moment–rotation response become less pronounced as \( D/B \) increases. In both cases, the moment–rotation loops reveal a great potential for the soil–foundation system to dissipate a large amount of energy during cyclic loading.

(e) Remarkable differences, however, are noticed with respect to the cyclically accumulating settlements as a function of \( \text{FSV} \). The reduced stresses transmitted onto the base soil would be expected to produce less settlement than for surface foundations. In fact, the lightly loaded system ended up with negligible residual settlement for both \( D/B = 0.5 \) and \( D/B = 1 \), after the
first few cycles during which the footing alternately uplifted and settled. On the other hand, the heavily loaded system displayed a substantial residual settlement of about 0.04B for D/B=0.5 and 0.06B for D/B=1, as its response is characterised by minimal uplifting and excessive soil yielding for the larger angles of rotation.

The effect of loading amplitude on settlement
This is further investigated by imposing five cycles of constant rotation amplitude, $\theta = 0.0005, 0.0025, 0.02, 0.05$ and 0.1 rad. Four safety factors, $FSV = 14, 5, 4$ and 3, are considered, while $D/B = 1$. Thus, a wide range of responses is covered, from nearly elastic to strongly inelastic. Figure 22 shows the residual settlement plotted against the imposed rotation amplitude. At relatively small angles of rotation, increasing amplitude implies linearly increasing settlement in all but the very large FSV case. At larger rotation amplitudes, the rate of settlement increase diminishes, until eventually settlements start levelling off or even decreasing. As a result, at large angles of rotation the lightly loaded systems end up with residual uplift, while the heavily loaded systems end up with reduced settlement. The ‘turning point’ is a critical angle of rotation, which appears to increase as the factor of safety decreases (for FSV = 3 this critical angle has not been reached yet at 0.1 rad.)

SUMMARY AND CONCLUSIONS
The static response of embedded square-in-plan foundations under undrained conditions has been studied with a 3D FE model. Soil (material) inelasticity and soil–foundation interface (geometric and material) non-linearity have been properly taken into account. Emphasis was given to the ultimate capacity under combined lateral (horizontal and vertical) loading, but the role of the vertical force was also addressed. Simplified models were developed with the sole purpose of elucidating the phenomena observed numerically. Some attention has been given to the effect of embedment on soil stiffnesses in pure horizontal translation and pure rotation.

For the (slow) cyclic loading the focus has been on the rotational response of embedded foundations carrying a simple slender structure. Numerical moment–rotation $(M-\theta)$ and settlement–rotation $(w-\theta)$ loops have been compared with the monotonic curves of the system, with particular emphasis on the importance of the static factor of safety $FSV$ against vertical ultimate capacity of the foundation. The rotation exceeded substantially the angle of maximum moment capacity of the foundation, and $P-\Delta$ effects played their role in function of FSV and the depth of embedment.

An interesting finding, which had not been noted explicitly in the literature before, was that the effect of embedment on increasing bearing capacity is much larger than the effect on increasing stiffness: almost four times larger for the horizontal motion and two times larger for rotation. Physical interpretations of these differences have been outlined.

At the soil–foundation interfaces detachment (with subsequent gap formation) and slippage (with shear tractions limited to adhesion) have been shown to exert a rather dramatic effect in reducing the size and even changing the shape of the failure envelopes. The relative magnitude of the vertical load affects mainly the uplifting of the base; its effect on the overall ultimate performance is limited.

Along the $QM$ failure envelope of a fully bonded embedded foundation the mechanisms of failure change significantly from point to point. For the four characteristic points, $M_{Q=14}$, $M_{Q=5}$, $M_{Q=2}$ and $M_{Q=4}$, an and $M_{Q=1}$, the observed mechanisms are, respectively: a clock-wise rotation about a point nearly $2/3$ of $D$ from the top ($M_{Q=14}$); a mixed mode consisting of rotation below the base about its centre and of passive- and active-type wedges along the sides normal to the direction of loading ($M_{Q=5}$); two clear wedges, a passive and an active on each of the two sides, with simple slippage along the base ($Q_{Q=2}$); and a counter-clock-wise rotation about a point at the top, like a pendulum ($Q_{Q=4}$). Of course, the two sidewalks that are parallel to the loading plane experience shear failure in all cases; that may involve mainly torsion, or horizontal and vertical translation, or combinations of these. With the non-linear (tensionless and sliding) interfaces, the mechanisms for the above four loading combinations change mainly due to the detachment that takes place; and the previous symmetry is lost.

The cyclic response of embedded foundations is summarised in the form of the foundation moment–rotation $(M-\theta)$ and settlement–rotation $(w-\theta)$ loops of a rotationally loaded system. The former loops are well enveloped by the corresponding monotonic curves. The latter loops show a small or large accumulation of settlement. The static vertical factor of safety $FSV$ controls the nature of these loops, perhaps as much as (and in some cases even more than) it affects the value of the ultimate resistance. With high $FSV$, interface (geometric) non-linearities prevail, with only minor inelastic (material) action in the soil and hence a small accumulation of settlement. The opposite is true when $FSV$ is small.

Closed-form expressions for the presented failure envelopes (although with the top of the foundation as the reference point) have been given by Ntritisos (2012) and recently by Karapiperis & Gerolymos (2014) and Zafeirakos (2014).

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NOTATION
- $A_{sidewall}$ area of foundation sidewall (= DB)
- $B$ width of footing
- $C$ initial kinematic hardening modulus
- $c_a$ cohesion at soil–foundation interface (= $aS_a$)
REFERENCES


